# Theoretical study of the cyclization of $\alpha$-iminothioaldehydes into dihydrothiazoles 

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#### Abstract

When synthesized by retro-Diels-Alder reaction under flash vacuum thermolysis conditions, $\alpha$ iminothioaldehydes undergo a rearrangement to 2,3 -dihydro-1,3-thiazoles. These secondary enamines in turn are transformed into the more stable 2,5-dihydro-1,3-thiazoles upon standing in $\mathrm{CDCl}_{3}$ solution. The route from $\alpha$-iminothioaldehydes to dihydro-1,3-thiazoles has been studied by ab initio molecular orbital theory. Geometries of stationary points on the [ $\mathrm{C}_{3} \mathrm{H}_{5} \mathrm{NS}$ ] potential energy surface were optimized using HF/6-31G(d,p) and MP2/6-3I(d,p) calculations. Relative energies were estimated using MP4STDQ [and in some cases QCISD(T)] calculations and the 6-31 $+\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set and corrected for zero-point vibrational energies. Intrinsic reaction coordinate calculations at $\mathrm{HF} / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level were performed in order to check the profile of the proposed rearrangement. These calculations support the experimental findings. The most energy requiring step was found to be trans $\rightarrow$ cis isomerization of the methylimino group $\left[27.8 \mathrm{kcal} \mathrm{mol}^{-1}\right.$ at the $\operatorname{QCISD}(\mathrm{T}) / 6-31+\mathrm{G}(\mathrm{d}, \mathrm{p})$ level]. $\dagger$


1,3-Thiazoles are the most widely studied $\mathrm{S}, \mathrm{N}$-heterocycles. ${ }^{1}$ Among their dihydro derivatives, the more easily available are those of the 4,5-dihydro series. ${ }^{2}$ The other series, 2,3- and 2,5dihydro, are little known. ${ }^{3}$ In particular, 2,3-dihydro-1,3thiazoles, without a substituent on the nitrogen atom, were unknown before this work. ${ }^{4}$ They are examples of secondary enamines and, as such, can be expected to be readily transformed into their more stable imino tautomers, 2,5-dihydro-1,3-thiazoles. Such isomerization is well documented for simple primary and secondary vinyl amines. ${ }^{5}$

A simple retro-synthetic analysis of 2,3-dihydro-1,3-thiazole (1a) (Scheme 1) shows that it could be obtained by a C(2)-S

connection from the enethiol 2a. This enethiol can be regarded either as the tautomer of $N$-(2-thioxoethyl)methanimine or as the product resulting from a $[1,5]$-sigmatropic rearrangement of the $\alpha$-iminothioaldehyde 3a. Indeed, when in the course of our work on thioaldehydes' reactivity, ${ }^{6}$ we tried to synthesize compound 3a, we observed its cyclization into 1a. ${ }^{4}$ The aim of this paper is to discuss in detail the mechanism leading from $\alpha$ iminothioaldehydes to dihydrothiazoles.

## Synthesis and thermolysis of imines 4a, b

Flash vacuum thermolysis (FVT) has proved to be an efficient technique for the generation and the spectroscopic observation of transient species. ${ }^{7}$ Among the reactions that can be used to synthesize reactive thiocarbonyl compounds under FVT

[^0]conditions, the retro-Diels-Alder reaction ${ }^{8}$ is one of the most powerful. ${ }^{6}$
The precursors $\mathbf{4 a}, \mathbf{b}$ (Scheme 2) were prepared in two steps


Scheme 2 Reagents and conditions: i, $\mathrm{Et}_{3} \mathrm{~N}, \mathrm{CHCl}_{3}$; ii, DIBAL-H, PhMe; iii, $\mathrm{R}_{2} \mathrm{CHNH}_{2}, \mathrm{CH}_{2} \mathrm{Cl}_{2}$; iv, FVT, $500^{\circ} \mathrm{C}$, -DMA
from nitrile 5. This nitrile was obtained by trapping thioformyl cyanide 6, generated in situ by the base induced elimination of HCN from cyanomethylthiocyanate, with 9,10-dimethylanthracene (DMA). ${ }^{9}$ Compound 5 was then reduced by DIBAL-H (diisobutylaluminium hydride) to give the corresponding aldehyde, which upon treatment with methyl- or isopropyl-amine gave the two imines $4 \mathbf{a}, \mathbf{b}$.

The thermolysis of these imines was effected at $500^{\circ} \mathrm{C} .9,10$ Dimethylanthracene condensed immediately at the oven exit, while the obtained volatile products were trapped on a cold finger (liquid $\mathrm{N}_{2}$ ) coated with a $\mathrm{CDCl}_{3}-\mathrm{CFCl}_{3}$ mixture for NMR analysis. The NMR spectra were recorded at $-60^{\circ} \mathrm{C}$. The obtained data, summarized in Scheme 3, clearly demonstrated the formation of the 2,3 -dihydro-1,3-thiazoles 1a, b. The corresponding 2,5 -dihydro-1,3-thiazoles $7 \mathbf{7 a}, \mathbf{b}^{10}$ were


1a -60 to $20^{\circ} \mathrm{C}$


79


1b



7b

Scheme $3{ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}$ NMR chemical shifts $\delta$ and coupling constants $J_{\mathrm{H}-\mathrm{H}}$ of compounds $\mathbf{1 a}, \mathbf{b}$ and $\mathbf{7 a}, \mathbf{b}$
also detected in the spectra (at $-60^{\circ} \mathrm{C}: \mathbf{1 a}: 7 \mathbf{a}=8: \mathbf{1 ;} \mathbf{1 b}: 7 \mathbf{b}=$ 4:1). Upon warming the NMR solution, the secondary enamines $1 \mathbf{a}, \mathbf{b}$ disappeared and at room temperature the imines $7 a, b$ were observed alone.

## Calculations

The calculations were carried out on an IBM RS6000/340 workstation with the GAUSSIAN92 quantum chemical program. ${ }^{11}$ Geometries of the stationary points were fully optimized without symmetry constraints at the single determinantal Hartree-Fock (HF) level of theory using the 6$31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set. ${ }^{12}$ Harmonic vibrational frequencies were computed in order to characterize stationary points (minima and first-order saddle points, the latter having exactly one imaginary frequency) and to estimate the zero-point vibrational energy (ZPVE, scaling factor of 0.89 ) contributions to the relative energies between stationary points.

After transition states were obtained, intrinsic reaction coordinate (IRC) ${ }^{13}$ calculations were carried out to follow the reaction path in both directions, leading back to the reactant as to the product. In these calculations, the geometry is optimized at each point along the reaction path as the system decreases in energy to verify that the saddle points obtained were associated with the reaction path originally assumed.

All structures were further optimized at the second order Moller-Plesset perturbation theory [MP2(fc)/6-31G(d,p)] level. Relative energies between stationary points were estimated using total energies computed at the fourth order perturbation theory (MP4SDTQ) and, in some cases, quadratic configuration interaction [QCISD(T)] levels with the 6-31 + $\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set ${ }^{12}$ using MP2 optimized geometries. The total energies in Hartree of the reference compound 3a (trans,s-trans) at the various levels of calculation are as follows: - 568.43820 $[\mathrm{HF} / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})], \quad-569.13320 \quad[\mathrm{MP} 2 / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})]$, -569.14551 [MP2/6-31G(d,p)], - 569.22136 [MP4SDTQ/6 $31+\mathrm{G}(\mathrm{d}, \mathrm{p})]$ and $-569.22312[\mathrm{QCISD}(\mathrm{T}) / 6-31+\mathrm{G}(\mathrm{d}, \mathrm{p})]$.

NBO analyses were performed with the program NBO version 3.1 which was built into link 607 of the GAUSSIAN92 program. ${ }^{14}$ Quantitative analysis of the hyperconjugation
interactions between various NBOs was obtained using the NBO Fock matrix deletion procedure. ${ }^{15 a}$ The deletion procedure is not self-consistent, but as long as the particular interactions in the Fock matrix that have been set to zero are not strongly coupled with other interactions, the error in the energy is not significant. ${ }^{15 a}$ Also, the energies of individual orbital interactions are not strictly additive. ${ }^{15 b}$

## Results and discussion

## Conformers of the $\alpha$-iminothioaldehyde 3a and their interconversions

Fig. 1 lists selected HF and MP2 parameters of local minima and transition structures involved in this process. The corresponding zero-point vibrational and relative energies are summarized in Table 1. Equilibrium geometries of the s-trans isomers are reached for $C_{\mathrm{s}}$ symmetry, while a non-planar structure is obtained for the s-cis rotamers. A $\angle \mathrm{SCCN}$ dihedral angle of $18.3^{\circ}(\mathrm{HF})$ and $14.2^{\circ}(\mathrm{MP} 2)$ is calculated for the trans,scis isomer and of $44^{\circ}(\mathrm{HF})$ and $29.4^{\circ}$ (MP2) for the cis,s-cis isomer. The HF and MP2 optimized geometries reported in Fig. 1 show that the absolute values of bond lengths depend on the level of calculation. As it can be seen, inclusion of electron correlation causes a general lengthening of the bond distances, with the exception of the $\mathrm{C}-\mathrm{C}$ bond which is shortened at the MP2/6-31G(d,p) level by $c a .0 .01 \AA$.

Another structural feature of the stable isomers of $\mathbf{3 a}$ is the conformational preference of the methyl group. HF/6-31G(d,p) calculations indicate that the lowest energy conformation has a $\mathrm{C}-\mathrm{H}$ bond lying in the plane of the $\mathrm{C}-\mathrm{N}-\mathrm{C}\left(\mathrm{H}_{3}\right)$ group and eclipsing the $\mathrm{C}=\mathrm{N}$ bond, while the staggered conformation was characterized as the transition state between two equivalent eclipsed rotamers. The rotational barrier is rather large ( 2.1 kcal $\mathrm{mol}^{-1}$ ) in the case of the two s-trans isomers. This tendency of the methyl group to eclipse the $\mathrm{C}=\mathrm{N}$ bond would be attributed to the nitrogen lone pair $\mathrm{n}_{\mathrm{N}} \rightarrow$ antibond $\sigma^{*}{ }_{\mathrm{CH}}$ orbital interaction which would be maximized for the eclipsed structure. If this interpretation is correct, one would expect that the substitution of the $\mathrm{C}-\mathrm{H}$ eclipsed bond by a $\mathrm{C}-\mathrm{F}$ bond, which enhances the hyperconjugative interaction, will stabilize the eclipsed rotamer


Fig. 1 HF/6-3IG(d,p) and MP2/6-31G(d,p) (italic) optimized geometrical parameters of the equilibrium and transition structures for the 3a(trans,strans $\rightarrow \mathbf{3 a}($ cis,s-cis) isomerization: bond lengths in $\AA$, angles in degrees. (a) 3a(trans,s-trans), (b) 3a(trans,s-cis), (c) 3a(cis,s-trans), (d) 3a(cis,s-cis), (e) TS1 (161i), ( $f$ ) TS1' (467i), (g) TS2 (432i), (h) TS2' (90i).

Table 1 Zero-point energies (ZPVE) ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) and relative energies ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) of the $\left(\mathrm{C}_{3} \mathrm{H}_{5} \mathrm{NS}\right)$ stationary points considered at different levels of theory

| Point | $\begin{aligned} & \text { ZPVE }{ }^{\text {a }} \\ & \text { energy } \end{aligned}$ | Relative energy |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $\mathrm{HF}^{\text {b }}$ | $\mathrm{MP}^{\text {c }}$ | MP2 $2^{\text {c.d }}$ | MP4SDTQ ${ }^{\text {c.d }}$ | QISD(T) ${ }^{\text {c.d }}$ | MP4 + ZPVE ${ }^{\text {e }}$ |
| 3a(trans,s-trans) | 45.6 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 |
| 3a(trans,s-cis) | 45.4 | 6.4 | 5.2 | 5.3 | 5.0 | 6.2 | 4.8 |
| 3a(cis,s-trans) | 45.7 | 6.5 | 5.3 | 5.6 | 5.4 | - | 5.5 |
| 3a(cis,s-cis) | 45.6 | 10.3 | 8.2 | 8.7 | 8.2 | 8.6 | 8.2 |
| 2a(gauche) | 44.1 | 8.3 | 9.9 | 8.8 | 9.4 | - | 7.9 |
| 2a(s-trans) | 44.2 | 4.4 | 7.1 | 6.4 | 6.8 | - | 5.4 |
| 2a(s-cis) | 44.3 | 2.7 | 4.9 | 4.0 | 4.6 | - | 3.3 |
| 8a(s-trans) | 46.5 | 18.7 | 13.3 | 12.3 | 14.4 | - | 15.3 |
| 8a(s-cis) | 46.8 | 22.5 | 17.1 | 15.7 | 17.7 | -- | 18.9 |
| $1 \mathbf{1 a}$ | 47.3 | $-3.2$ | $-4.7$ | -6.4 | $-3.8$ | - | $-2.1$ |
| 7 a | 47.1 | $-15.5$ | $-16.7$ | $-16.2$ | $-14.2$ | - | - 12.7 |
| TS1 | 45.1 | 8.5 | 8.1 | 7.5 | 7.1 | - | 6.6 |
| TS1 ${ }^{\prime}$ | 44.4 | 31.0 | 32.0 | 32.3 | 32.5 | 32.6 | 31.3 |
| TS2 | 44.3 | 34.1 | 34.3 | 34.8 | 34.8 | 35.0 | 33.5 |
| TS2 ${ }^{\prime}$ | 45.2 | 10.9 | 9.8 | 9.6 | 9.2 | - | 8.8 |
| TS3 | 42.8 | 54.7 | 31.7 | 30.8 | 35.0 | 37.1 | 32.2 |
| TS4 | 43.8 | 8.9 | 11.5 | 10.3 | 10.8 | - | 9.0 |
| TS5 | 43.9 | 7.4 | 11.4 | 10.7 | 11.0 | -- | 9.3 |
| TS6 | 43.0 | 29.4 | 19.5 | 19.0 | 21.8 | - | 19.2 |
| TS7 | 46.2 | 32.0 | 37.6 | 34.0 | 35.3 | - | 35.9 |
| TS8 | 46.4 | 26.5 | 21.8 | 19.9 | 21.6 | - | 22.4 |
| TS9 | 43.3 | 98.3 | 71.1 | 71.7 | 72.4 | - | 70.1 |

${ }^{a}$ From HF/6-31(d,p) harmonic vibrational wavenumbers and scaled by $0.89 .{ }^{b}$ Based on HF/6-31G(d,p) geometries. ${ }^{\text {c Based on MP2(fc)/6- }}$ $31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ geometries. ${ }^{d}$ With the $6-31+\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set; core orbitals are frozen. ${ }^{e}$ Including MP4 relative energies and zero-point vibrational energies.

Table 2 Energy contributions of $n_{N} \rightarrow \sigma^{*}{ }_{C X}$ delocalization and energy differences between eclipsed and staggered conformations of $\mathrm{S}(\mathrm{CH})_{2} \mathrm{NCH}_{2} \mathrm{X}, \mathrm{X}=\mathrm{H}$ and $\mathrm{F}^{a}$

|  | $\mathrm{X}=\mathrm{H}$ |  | $\mathrm{X}=\mathrm{F}$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Eclipsed | Staggered | Eclipsed | Staggered |
| $\mathrm{n}_{\mathrm{N}} \rightarrow \sigma^{*}{ }^{\text {CH }}{ }^{6}$ | -0.4 | -2.2 | -0.4 | -1.8 |
| $\mathrm{n}_{\mathrm{N}} \rightarrow \sigma^{*}{ }_{\mathrm{CX}}$ | -7.7 | -2.6 | $-13.3$ | -8.7 |
| $\Sigma^{c}$ | -8.5 | -7.1 | - 14.0 | $-12.7$ |
| $\Delta E_{\text {del }}{ }^{\text {e }}$ | 0.0 | 4.8 | 0.0 | 8.5 |
| $\Delta E_{\text {Lew }}{ }^{\text {d }}$ | 0.0 | -2.8 | 0.0 | -3.5 |
| $\Delta E$ | 0.0 | 2.0 | 0.0 | 5.0 |

${ }^{a}$ Energies are in $\mathrm{kcal} \mathrm{mol}^{-1} \quad{ }^{b}$ There are two identical orbit interactions of this type, the CH bonds lying above and below the $\mathrm{C}=\mathrm{N}-\mathrm{C}\left(\mathrm{CH}_{2} \mathrm{X}\right)$ plane. ${ }^{c}$ Total of the energies for the individual lone pair to vicinal $\mathrm{CH}(\mathrm{X})$ antibond delocalizations. ${ }^{d}$ Difference in the energies of the Lewis structure obtained by setting all orbitals' interactions to zero. ${ }^{e}$ Difference in delocalization energies ( $\Delta E_{\mathrm{del}}=$ $\left.\Delta E_{\text {tot }}-\Delta E_{\text {Lew }}\right)$.
and increase the rotational barrier. Calculations have been performed for the trans,s-trans fluorine derivative $\mathrm{S}=\mathrm{CH}-\mathrm{CH}=$ $\mathrm{N}-\mathrm{CH}_{2} \mathrm{~F}$. There is a total of four different structures corresponding to stationary points on the $\mathrm{HF} / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ potential energy surface. As expected, the $\mathrm{C}-\mathrm{F}$ and $\mathrm{C}=\mathrm{N}$ bonds are eclipsed in the most stable conformation. In the other stable structure a $\mathrm{C}-\mathrm{H}$ bond eclipses the $\mathrm{C}=\mathrm{N}$ one. The two staggered rotamers ( $\mathrm{C}-\mathrm{F}$ or $\mathrm{C}-\mathrm{H}$ and $\mathrm{C}=\mathrm{N}$ anti) are identified as firstorder saddle points. The energy difference between the eclipsed (CF) and staggered (CF) conformations is equal to 5.0 kcal $\mathrm{mol}^{-1}$. In order to definitively check our hypothesis, we have carried out a Fock matrix analysis with the NBO program ${ }^{14}$ for the trans, s-trans $\mathrm{S}=\mathrm{CH}-\mathrm{CH}=\mathrm{N}-\mathrm{CH}_{2} \mathrm{X}(\mathrm{X}=\mathrm{H}, \mathrm{F})$ compounds. Orbital interactions of interest (nitrogen lone pair and antibond $\sigma_{c x}^{*}$ orbitals of $\mathrm{CH}_{2} \mathrm{X}$ group) were calculated by setting the corresponding off-diagonal Fock matrix elements to zero. The results are summarized in Table 2. As predicted, the NBO
analysis indicates the importance of $\mathrm{n}_{\mathrm{N}} \rightarrow \sigma_{\mathrm{CX}}^{*}$ hyperconjugation when $C-X$ and $C=N$ are eclipsed and this contribution increases significantly for $X=F$. However, a close examination of Table 1 shows that the variation of the sum $\Sigma$ of these interactions is almost constant whatever X and does not parallel $\Delta E_{\text {del }}$. Hence, the energy difference between eclipsed and staggered structures cannot be rationalized quantitatively by $\mathrm{n}_{\mathrm{N}} \rightarrow \sigma^{*}{ }_{\mathrm{CX}}$ orbital interactions alone.

The most stable 3a isomer is the trans,s-trans one (see Table 1) and the order of stability is the same whatever the level of calculations. There are two possible pathways from $\mathbf{3 a}$ (trans,strans) to $\mathbf{3 a}(c i s, s-c i s)$ (Fig. 2). The first involves a rotation about the $\mathrm{C}-\mathrm{C}$ bond giving the trans,s-cis rotamer via the transition state TS1 (Fig. 1) followed by inversion at the nitrogen atom forming the cis,s-cis compound via TS2. In the second pathway, inversion at nitrogen occurs first via TS1' and is followed by rotation about the C-C bond via TS2'. TS1 and TS2' are reached for a $\angle S C C N$ dihedral angle of about $90^{\circ}$; only TS2 has the $C_{s}$ symmetry; the $\mathrm{CNC}\left(\mathrm{H}_{3}\right)$ bond angle is close to $180^{\circ}$ in both TS1' and TS2. Table 3 reports the energy barriers of these steps. It appears that (i) barrier levels are reasonably well predicted at the HF level, the correlation corrections make little differences; this result can be related to the fact that, in the corresponding transition structures, no bond is partially broken or formed. (ii) The barrier heights for the inversion at nitrogen are, as expected, substantial (ca. 30 $\mathrm{kcal} \mathrm{mol}^{-1}$ ); similar results were obtained by Pople et al. in their study of the potential energy surface of $\mathrm{H}_{2} \mathrm{C}=\mathrm{NH} .{ }^{16}$ (iii) At the best level of calculation [ZPE/QCISD $(\mathrm{T}) / 6-31+\mathrm{G}(\mathrm{d}, \mathrm{p})$ ] the inversion at the nitrogen of the trans,s-trans isomer constitutes the most energy-demanding step.

## From 3a(cis,s-cis) to the enethiol 2a

Fig. 3 displays selected HF and MP2 parameters of stationary points relevant to this process. TS3 is the transition structure for the [1,5]-sigmatropic hydrogen shift connecting 3a and 2a. IRC calculations at the 6-31G(d,p) level revealed that this transition state connects exactly the $\mathbf{3 a}(c i s, \mathrm{~s}-c i s)$ and $\mathbf{2 a}($ gauche $)$ isomers.

Table 3 Energy barriers $\Delta E^{\ddagger}\left(\mathrm{kcal} \mathrm{mol}^{-1}\right)$ and activation energy $\left(\mathrm{kcal} \mathrm{mol}^{-1}\right)$ of the steps involved in the $\alpha$-iminothioaldehyde rearrangement

|  | $\Delta E^{\ddagger}$ |  |  |  |  | Activation energy ${ }^{d}$ MP4(QCISD) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{HF}^{\text {a }}$ | MP2 ${ }^{\text {b }}$ | MP2 ${ }^{\text {b.c }}$ | MP4SDTQ ${ }^{\text {b,c }}$ | QCISD( ) $^{\text {b.c }}$ |  |
| TS1 | 8.5 | 8.1 | 7.5 | 7.1 | - | 6.6 |
| TS1 ${ }^{\prime}$ | 31.0 | 32.0 | 32.3 | 32.5 | 32.6 | 31.3 (31.4) |
| TS2 | 27.7 | 29.1 | 29.5 | 29.7 | 28.8 | 28.7 (27.8) |
| TS2 ${ }^{\prime}$ | 4.4 | 4.6 | 4.4 | 4.1 | - | 3.3 |
| TS3 | 44.4 | 23.5 | 22.0 | 26.8 | 28.5 | 24.0 (25.7) |
| TS4 | 1.4 | 1.6 | 1.6 | 1.4 | - | 1.5 |
| TS5 | 3.0 | 4.4 | 4.3 | 4.2 | - | 3.9 |
| TS6 | 26.6 | 14.6 | 15.0 | 17.1 | - | 15.9 |
| TS7 | 13.3 | 24.4 | 21.7 | 20.9 | - | 20.6 |
| TS8 | 4.0 | 4.7 | 4.2 | 4.0 | - | 3.5 |
| TS9 | 101.5 | 71.1 | 71.7 | 72.4 | - | 72.2 |

${ }^{a}$ Based on HF/6-3IG(d,p) geometries. ${ }^{b}$ Based on MP2/6-3IG(d,p) geometries. ${ }^{c}$ With the $6-31+\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set; core orbitals are frozen.
${ }^{d}$ Including zero-point vibrational energies.


Fig. 2 Theoretical estimates of the relative energies at the MP4/6$31+\mathrm{G}(\mathrm{d}, \mathrm{p})$ level for the whole process from $\mathbf{3 a}($ trans, $\mathrm{s}-$ trans $)$ to $\mathbf{8 a}$

The main difference between HF and MP2 geometries concern the lengths of forming and breaking bonds which are $0.05 \AA$ shorter at the MP2 level. Therefore, the transition state is predicted to be more tightened at this level. Pauling bond orders ${ }^{17} \mathrm{n}_{\mathrm{p}}$ provide a useful measure of the extent of bond making ( bm ) and bond breaking ( bb ) in the transition structures, $\operatorname{Ln}\left(n_{p} / n_{0}\right)=\left(d_{0}-d\right) / 0.3$, where $d$ and $d_{0}$ are the $\mathrm{C}\left(\mathrm{H}_{2}\right)$ H bond lengths respectively in TS3 and $\mathbf{3 a}(\mathrm{cis}, \mathrm{s}-\mathrm{cis})$ and $\mathrm{S}-\mathrm{H}$ bond lengths in TS3 and product for bond making; $n_{0}$ is equal to 1 for single bonds. $n_{\mathrm{p}}(\mathrm{bb})$ Values of 0.346 (HF) and 0.426 (MP2) are calculated. The $n_{\mathrm{p}}(\mathrm{bm})$ values are lower, slightly at the HF level ( 0.327 ), more at the MP2 level ( 0.332 ). Bond index values are not decisive to determine the 'lateness' of TS3.

Further information concerning TS3 can be obtained using the natural bond analysis (NBO) and natural localized molecular orbital analysis (NLMO) methods. ${ }^{18}$ The principal resonance structure possesses a covalent $\mathrm{S}-\mathrm{H}$ bond and the NLMO bond orders, ${ }^{19}$ which give a covalent contribution to bonding, are equal to 0.512 for the bond between the sulfur atom and the migrating H and only 0.291 for the bond between this H and the carbon atom. It can be concluded that TS3 resembles the product more than the reactant. This conclusion is further supported by the IRC calculations. From the values of the net reaction coordinate ( $s=0$ corresponds to the transition structure) for the cis,s-cis- $\alpha$-iminothioaldehyde $(s=9.05)$ and the gauche-enethiol ( $s=-7.55$ ) it becomes apparent that TS3 is structurally closer to the enethiol than to the thioaldehyde. At the $\mathrm{HF} / 6-3 \mathrm{l} G(\mathrm{~d}, \mathrm{p})$ level, $\mathbf{2 a}($ gauche $)$ lies $2.0{\mathrm{kcal} \mathrm{mol}^{-1} \text { below }}^{\text {a }}$
$\mathbf{3 a}(c i s, s-c i s)$ in energy. When correlation is introduced, a reverse order is obtained. However, when ZPVE correction is taken into account, this rearrangement is found to be nearly athermic ( $-0.3 \mathrm{kcal} \mathrm{mol}^{-1}$ at the MP4SDTQ + ZPVE level).

A barrier of $44.4 \mathrm{kcal} \mathrm{mol}^{-1}$ was calculated for this step at the HF/6-13G(d,p) level. As expected, when correlated methods which provide a more realistic description of the bond-breaking and -making processes were used, the barrier was significantly lowered (see Table 3). It is noteworthy that, at each level of theory, the barrier is substantially lower than that calculated for the [1,5]-hydrogen shift in ( $Z$ )-penta-1,3-diene (which varies in the range 36.5 (MP2)-58.1 (HF) $\mathrm{kcal} \mathrm{mol}^{-1}$ ). ${ }^{20}$ At the higher level of theory, the activation energy was calculated to be 25.7 $\mathrm{kcal} \mathrm{mol}^{-1}$, lower than that of the limiting $\mathbf{3 a}($ trans,s-trans) $\longrightarrow \mathbf{3 a}($ cis,s-trans $)$ step ( $31.4 \mathrm{kcal} \mathrm{mol}^{-1}$ ). This result is in agreement with an expected rapid [1,5]-sigmatropy.

As mentioned previously, IRC calculations indicate that the product resulting from the [1,5]-sigmatropy is the gaucheenethiol. We have located two other stable conformers of $\mathbf{2 a}$ each having a $C_{\mathrm{s}}$ symmetry (Fig. 3). The s-trans rotamer is predicted to be more stable than the gauche one by 2.5 kcal $\mathrm{mol}^{-1}$ (MP4 + ZPVE level) and the absolute minima of $\mathbf{2 a}$ is the s-cis rotamer which lies $2.1 \mathrm{kcal} \mathrm{mol}^{-1}$ below the s-trans rotamer. The $\mathbf{2 a}($ gauche $) \longrightarrow \mathbf{2} \mathbf{a}(s-t r a n s)$ rotational barrier (TS4) is low (ca. $1.5 \mathrm{kcal} \mathrm{mol}^{-1}$ at both levels of calculation). TS4 is reached by rotation around both $\mathrm{C}-\mathrm{S}$ and $\mathrm{C}-\mathrm{N}$ bonds. Thus in $\mathbf{2 a}($ gauche $)$, the $\angle \mathrm{HSCC}$ and $\angle \mathrm{CNCC}$ dihedral angles are $141.4^{\circ}$ and $53.4^{\circ}$, respectively (MP2). These values become $165.9^{\circ}$ and $99.0^{\circ}$ (MP2) in TS4. 2a(s-cis) is obtained via TS5. This step proceeds with an activation barrier of about $4.0 \mathrm{kcal} \mathrm{mol}^{-1}$. Consequently, the whole transformation $\mathbf{2 a}($ gauche $) \longrightarrow \mathbf{2 a}(s-c i s)$ is a facile step.

It is of interest to examine the origin of the largest stability of the s-cis isomer. The presence of an intramolecular $\mathrm{N} \cdot \mathrm{H}-\mathrm{S}$ hydrogen bond seems to be a natural explanation and, for the scis isomer, a stronger nitrogen lone pair $n_{N}-\sigma_{S H}^{*}$ interaction is expected. ${ }^{18}$ Table 4 reports the main energy terms resulting from a NBO deletion procedure applied to the $\mathbf{2 a}$ stable isomers. Surprisingly, the $n_{N} \rightarrow \sigma^{*}$ sH is weak and its variation is not sufficient to explain the relative stabilities of $\mathbf{2 a}$ isomers. Our results show that the whole delocalization energy $E_{\text {del }}$ is the factor determining the stability of these rotamers.

## From the enethiol 2a to the zwitterion 8a

The transition structure TS6 (Fig. 4) describes the migration of the H atom from $\mathrm{S}\left(2-\mathrm{H}\right.$ bond) to $\mathrm{N}\left(\mathrm{N}^{+}-\mathrm{H}\right.$ bond). At the two levels of geometry optimization, breaking and forming bond lengths are very close, as is the $\angle S H N$ bond angle. In TS6 the breaking $\mathrm{S}-\mathrm{H}$ bond is $0.325 \AA(\mathrm{HF}$ ) or $0.310 \AA$ (MP2) longer


Fig. $3 \mathrm{HF} / 6-3 \mathrm{IG}(\mathrm{d}, \mathrm{p})$ and MP2/6-31G(d,p) (italic) geometrical parameters for the 3a(cis,s-cis) $\longrightarrow \mathbf{2 a}$ transformation. (a) 2a(gauche), (b) $\mathbf{2 a}(\mathrm{gauche})$. (c) 2a(s-trans), (d) 2a(s-cis), (e) TS3 (2265i), ( $f$ ) TS3, ( $g$ ) TS4 (114i), (h) TS5 (226i).
than the corresponding bond in $\mathbf{2 a}(\mathrm{s}-\mathrm{cis})$. The forming $\mathrm{N}-\mathrm{H}$ bond is $0.274 \AA$ (HF) or $0.262 \AA$ (MP2) longer than that of the zwitterion 8a. The main feature of this transition state is its
planarity which requires that the orbital interactions are qualitatively different from those encountered in pericyclic reactions. Despite a large lengthening of the S-H bond in this


Fig. $4 \mathrm{HF} / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ and MP2/6-31(d,p) (italic) geometrical parameters for the $\mathbf{2 a} \longrightarrow \mathbf{8 a}$ rearrangement. (a) 8a(s-trans), (b) 8a(s-cis), (c) TS6 (1747i), (d) TS7 (191i).

Table 4 Contributions ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) of $\mathrm{n}_{\mathrm{N}} \rightarrow \sigma^{*}$ orbital interaction, delocalization as a total ( $\Delta E_{\text {del }}$ ) and Lewis energies ( $\Delta E_{\text {Lew }}$ ) to the energy differences ( $\Delta E_{\text {tot }}$ ) between the three $\mathbf{2 a}$ isomers

|  | $\mathbf{2 a}($ gauche $)$ | 2a(s-trans) | $\mathbf{2 a ( s - c i s )}$ |
| :--- | :---: | :---: | ---: |
| $\mathrm{n}_{N} \rightarrow \sigma^{*}{ }_{\text {sH }}$ |  |  |  |
| $E_{\text {del }}$ | 0.0 | -0.8 | -1.3 |
| $\Delta E_{\text {del }}$ | -226.5 | -227.8 | -223.1 |
| $\Delta E_{\text {Lew }}$ | 5.6 | 4.3 | 0.0 |
| $\Delta E_{\text {lot }}$ | -0.5 | -2.4 | 0.0 |

transition structure, NBO/NLMO analyses seem to indicate that TS6 is a rather early transition state: a S-H bond is present in the 'best' resonance structure and NLMO bond orders for $\mathrm{S}-\mathrm{H}$ and $\mathrm{N}-\mathrm{H}$ bonds are respectively 0.387 and 0.184 . However, the IRC which is depicted for this process in Fig. 5, gives values of the net reaction coordinate equal to -5.67 for the reactant 2a and to 3.08 for the product $8 \mathbf{8 a}$.
The IRC prediction of a transition state structurally close to 8a than to $\mathbf{2 a}$ can be related to the endothermicity of this process $\left(+12.0 \mathrm{kcal} \mathrm{mol}^{-1}\right.$ at the MP4 + ZPVE level) in accordance with the Hammond postulate. Thus, it appears that the difference between the geometric and electronic descriptions of bonds breaking and forming is stronger here.
The usual trends are observed in calculated energy barrier: MP2 underestimates it and MP4 begins to converge on a reasonable energy barrier (Table 3). The MP4(STDQ) + ZPVE energy of TS6 is $15.9 \mathrm{kcal} \mathrm{mol}^{-1}$ above $\mathbf{2 a ( s - c i s ) , ~ a ~ v a l u e ~}$ lower than the activation energy calculated for the $[1,5]$ sigmatropic shift.


Fig. 5 Potential energy profile for the $\mathbf{2 a} \longrightarrow \mathbf{8 a}$ rearrangement

Let us now comment on the structures of the intermediate 8a. As mentioned previously, IRC revealed that TS6 gives exactly the s-trans isomer. TS7 is the transition structure which connects $\mathbf{8 a}(\mathrm{s}-$ trans $)$ and $\mathbf{8 a ( s - c i s )}$ (see Fig. 4). The salient result emerging from the data listed in Fig. 4 is that geometrical parameters depend strongly on the level of calculation. The most conspicuous changes are the shortening of the $\mathrm{C}-\mathrm{N}$ and $\mathrm{C}-\mathrm{S}$ bonds and the lengthening of the $\mathrm{C}=\mathrm{C}$ and $\mathrm{N}=\mathrm{C}$ bonds when electron correlation corrections were included by MP2/6$31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ geometry optimizations. These modifications reveal changes in the $\pi$-character of these bonds and suggest that the contribution of the polar canonical form $\mathbf{B}$ (Scheme 4) as a resonance structure of $\mathbf{8 a}$ increases with correlation.
This analysis is supported by the NBO search of the Lewis


Scheme 4 Overall non-Lewis occupancies of the resonance structures A and $\mathbf{B}$ in $\mathbf{8 a}(\mathrm{s}-\mathrm{cis})$ and $\mathbf{8 a}(\mathrm{s}$-trans)
structures. In Scheme 4 are reported for each structure the nonLewis occupancies (the lower the non-Lewis density the better the corresponding structure is). When comparing these values, it can be concluded that, at the MP2 level, the participation of the resonance structure $\mathbf{B}$ to the description of $\mathbf{8 a}$ is more important. As a result, the correlated rotational barrier around the $\mathrm{C}-\mathrm{N}$ bond which lies in the range $20.9-24.4 \mathrm{kcal} \mathrm{mol}^{-1}$ is larger than that calculated at the HF level ( $13.3 \mathrm{kcal} \mathrm{mol}^{-1}$ ).

Whatever the level of calculation, the s-trans intermediate is found to be more stable than its s-cis rotamer (energy difference $=3.3-3.8 \mathrm{kcal} \mathrm{mol}^{-1}$ ). The NBO deletion procedure indicates that the former isomer is favoured by nonhyperconjugative effects, e.g. steric and electrostatic $\left[\Delta E_{\mathrm{T}}=3.5\right.$ $\mathrm{kcal} \mathrm{mol}{ }^{-1}, \Delta E_{\text {del }}=-1.1 \mathrm{kcal} \mathrm{mol}^{-1}, \Delta E_{\mathrm{Lew}}=4.6 \mathrm{kcal} \mathrm{mol}^{-1} \mathrm{at}$ the HF/6-31 $+\mathrm{G}(3, \mathrm{p})$ level with MP2 geometries].

From the zwitterion 8a to the cycles 1a and 7a
The geometries of the main points of these paths are presented in Fig. 6. The cyclization of $\mathbf{8 a}(\mathrm{s}$-cis) into $1 \mathbf{1 a}$ occurs via the transition state TS8. This ring closure is exothermic by 21.0 kcal $\mathrm{mol}^{-1}$. It is a facile step with an energy barrier of only 4.0 kcal $\mathrm{mol}^{-1}$. This value varies little with the level of calculation. TS8 is an early transition state with the $S \cdots C(5)$ length very extended $[2.818 \AA(\mathrm{HF})$ and $2.760 \AA$ (MP2)] and the $\mathrm{N}-\mathrm{C}(5)$ bond displaying essentially double bond character ( $1.337 \AA$ vs. $1.334 \AA$ in the reactant and $1.454 \AA$ in the product at the MP2 level). IRC calculations ( $s=5.16$ for the reactant $\mathbf{8 a}, s=6.45$ for the cyclic product 1a) as well as NBO/NLMO analysis [S . . C(5) bond order equal to 0.111 in TS8 and to 0.887 in 1a] support the geometrical predictions.

The ring closure leads to a conformer of 1 a having an axial $\mathrm{N}-\mathrm{H}$ bond (verified by the IRC calculation). The equatorial form is less stable than the axial one by $0.5 \mathrm{kcal} \mathrm{mol}^{-1}$ [MP4STDQ/6-31 + G(d,p) + ZPE//MP2/6-13G(d,p)].

Finally, we will consider the last step of this rearrangement, the [1,3]-sigmatropic $\mathrm{H}(10)$ shift, which experimentally occurs in $\mathrm{CDCl}_{3}$ solution. The symmetry forbidden suprafacial migration is expected to be an unfavourable pathway. The first order transition structures TS9 (Fig. 6) was located at the HF/6$31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level. However, the standard procedure at the MP2/6$31 G(d, p)$ level failed to converge on a transition structure. Due to the size of the system and prohibitive memory size requirements, it was unpractical to obtain MP2/6-31 (G(d,p) second derivatives and, instead, a directed grid search was used. The $\mathrm{N}-\mathrm{H}(10)$ and $\mathrm{C}(2)-\mathrm{H}(10)$ bonds were fixed and the other geometrical parameters were optimized. The forces of the fixed bonds and the total energy were used to direct the search. A conservative estimate would be that the bond lengths are optimized within the range $0.01-0.02 \AA$. There is a large discrepancy between some HF and MP2 bond lengths, but the two methods agree on a short $\mathrm{C}(1)-\mathrm{H}(10)$ distance in the transition structure [1.190 $\AA(\mathrm{HF}), 1.296 \AA$ (MP2)]. Consequently, it is not obvious that $\mathrm{H}(10)$ in TS9 transfers to N during the reaction. For this reason, we made use of the HF/6-31G(d,p) TS9 geometry to calculate the IRC. The distances between

Table 5 Distances (in $\AA$ ) between the migrating $\mathrm{H}(10)$ and $\mathrm{N}, \mathrm{C}(1)$ and $\mathrm{C}(2)$ for several points selected along the IRC for $1 \alpha \longrightarrow 7 a$ rearrangement

| Net reaction coordinate <br> $s$ | Distance $^{a}$ |  |  |
| :--- | :--- | :--- | :--- |
|  | $\mathrm{H}(10)-\mathrm{N}$ | $\mathrm{H}(10)-\mathrm{C}(1)$ | $\mathrm{H}(10)-\mathrm{C}(2)$ |
| $-2.94(\mathbf{1 a})$ | 1.002 | 2.009 | 2.937 |
| -1.60 | 1.057 | 1.491 | 2.140 |
| -0.60 | 1.304 | 1.261 | 1.834 |
| -0.20 | 1.464 | 1.203 | 1.697 |
| $0.00(\mathbf{T S} 9)$ | 1.549 | 1.190 | 1.630 |
| 0.20 | 1.635 | 1.189 | 1.566 |
| 0.40 | 1.720 | 1.200 | 1.503 |
| 0.80 | 1.885 | 1.253 | 1.379 |
| 1.20 | 2.039 | 1.343 | 1.263 |
| 2.80 | 2.610 | 1.861 | 1.091 |
| 4.00 | 2.948 | 2.081 | 1.085 |
| $4.49(7 \mathbf{7 a})$ | 2.980 | 2.071 | 1.082 |

" Atoms' numbering, see Fig. 6.
$H(10)$ and the $N, C(1)$ and $C(2)$ centres for various values of the net reaction coordinate are given in Table 5. It appears clearly that 1a and 7a are connected via TS9. It is also of interest to notice that in the range from $s=-0.2$ to $s=+0.4$, the $C(1)-$ $\mathrm{H}(10)$ distance remains almost constant while the other two distances vary noticeably in the meantime. The principal resonance structure found by the NBO search is shown in Scheme 5. It is consistent with the results given by the natural population analysis [NPA charges at $\mathrm{N}(3)$ and $\mathrm{H}(10)$ are respectively -0.749 and +0.532$]$. Calculated NLMO/NPA bond orders are $0.103[\mathrm{H}(10)-\mathrm{N}], 0.462[\mathrm{H}(10)-\mathrm{C}(1)]$ and 0.251 $[\mathrm{H}(10)-\mathrm{C}(2)]$. Again in this case variations of bond orders do not parallel closely geometrical trends [in view of elongation of bonds in TS9, a value of 0.251 for the $\mathrm{H}(10)-\mathrm{C}(2)$ bond is unexpected].


Scheme 5
The final product $7 \mathbf{a}$ is more stable than $\mathbf{1 a}$ by $10.6 \mathrm{kcal} \mathrm{mol}^{-1}$. Despite the exothermicity of this step, its activation energy was calculated to be large ( $72.2 \mathrm{kcal} \mathrm{mol}{ }^{-1}$ ). Thus, a thermal 1 a $\longrightarrow 7 \mathbf{a}$ rearrangement is very unlikely. This enamine $\longrightarrow$ imine rearrangement is most probably not a unimolecular process. ${ }^{5}$

## Conclusions

Fig. 2 summarizes the mechanism and energetics for the overall process. Apart from the last $\mathbf{1 a} \longrightarrow \mathbf{7 a}$ step, the highest


Fig. $6 \mathrm{HF} / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ and MP2/6-31G(d,p) (italic) geometrical parameters for the cyclization step. (a) 1a, (b) 1a, (c) 7a, (d) TS8 (195i), (e) TS9 (1935i), (f) TS9.
calculated barriers correspond to the 3a-trans $\longrightarrow$ cis isomerization and to the ( 1,5 )-sigmatropic shift leading to the enethiol $\mathbf{2 a}($ gauche $)$. At our best level of theory, i.e. $\operatorname{QCISD}(\mathrm{T}) / 6-31+$ $\mathrm{G}(\mathrm{d}, \mathrm{p}) / / \mathrm{MP} 2 / 6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})+$ ZPVE corrections, the activation energy of this second step was calculated to be lower ( 25.7 kcal $\mathrm{mol}^{-1}$ ) than that of the first one ( $27.8 \mathrm{kcal} \mathrm{mol}^{-1}$, see Table 3). Thus, the trans $\longrightarrow$ cis isomerization constitutes the most energy requiring step. Figs. 1, 3, 4 and 6 list the unscaled imaginary frequencies of all the transition structures. For TS2, the absolute value of the reaction coordinate frequency (431i) is much lower than that calculated for TS3 (2265i). To the extent that the Wigner and Bell methods for approximating tunnelling effect ${ }^{21}$ are a reasonable picture of physical reality, a larger tunnel correction is expected for the sigmatropic [1,5]hydrogen migration.

## Experimental

## 12-Formyl-9,10-dimethyl-9,10-dihydro-11-thia-9,10-ethano-

## anthracene

Compound $5^{9}(5 \mathrm{~g} ; 18 \mathrm{mmol})$ was dissolved in dry toluene ( 280 $\mathrm{cm}^{3}$ ) at $0^{\circ} \mathrm{C}$. A $1.5 \mathrm{~mol} \mathrm{dm}{ }^{-3}$ solution of DIBAL-H in toluene ( 54 mmol ) was then added dropwise. The reaction mixture was stirred from $0^{\circ} \mathrm{C}$ to room temperature for 12 h . After cooling back to $0^{\circ} \mathrm{C}$, the hydrolysis was performed with hydrochloric acid $\left(1 \mathrm{~mol} \mathrm{dm}^{-3}\right)$. The organic layer was separated and washed twice with water. Toluene was evaporated under reduced pressure and the resulting oil was submitted to column chromatography (eluent pentane-chloroform, $40: 60$ ). Yield $92 \%$; mp: $120-130^{\circ} \mathrm{C}$ (decomp.); $v\left(\mathrm{CDCl}_{3}\right) / \mathrm{cm}^{-1} 1700(\mathrm{CHO})$; $\delta_{\mathrm{H}}\left(\mathrm{CDCl}_{3} ; 300 \mathrm{MHz}\right) 2.16$ and $2.34(2 \mathrm{~s}, 6 \mathrm{H}, \mathrm{Me}), 3.50(\mathrm{~d}, 1 \mathrm{H}$,
$\mathrm{CH}), 7.2-7.4(\mathrm{~m}, 8 \mathrm{H}, \mathrm{Ar})$ and $8.42(\mathrm{~d}, 1 \mathrm{H}, \mathrm{CHO}) ; \delta_{\mathrm{C}}\left(\mathrm{CDCl}_{3} ; 75\right.$ $\mathrm{MHz}) 17.26$ and $18.29(2 \mathrm{Me}), 44.61$ and $46.86(\mathrm{C}), 63.44(\mathrm{CH})$, $118.91,119.28,121.90,122.82,126.26,126.85,127.11,140.00$, $142.63,144.73$ and 145.26 (Ar) and 195.98 (CHO); m/z 280 $\left(\mathrm{M}^{\cdot+}, 6 \%\right), 251\left(\mathrm{M}^{+}+\mathrm{CHO}, 1\right), 206$ (dimethylanthracene ${ }^{\bullet+}$, 100), 191 (methylanthracene ${ }^{+}, 60$ ), 178 (anthracene ${ }^{\cdot+}, 12$ ), 74 $\left(\mathrm{M}^{++}-\right.$DMA, 3) and 45 (10) (HRMS Calcd. $\mathrm{C}_{18} \mathrm{H}_{16} \mathrm{OS}$; 280.0922 Found: 280.0922).

## 12-( $N$-Methyliminomethyl)-9,10-dimethyl-9,10-dihydro-11-thia-9,10-ethanoanthracene (4a)

Gaseous methylamine was bubbled for 2 h at $0^{\circ} \mathrm{C}$ through a solution of 12-formyl-9,10-dimethyl-9,10-dihydro-11-thia-9,10ethanoanthracene ( 1 mmol ) in dichloromethane $\left(5 \mathrm{~cm}^{3}\right)$ containing molecular sieves ( $4 \AA$ ). The solvent was evaporated under reduced pressure. Compound $\mathbf{4 a}$ was obtained in quantitative yield. No aldehyde was detected in the NMR and IR spectra. Mp $125-130^{\circ} \mathrm{C}$ (decomp.); $v\left(\mathrm{CDCl}_{3}\right) / \mathrm{cm}^{-1} 1656$ $(\mathrm{C}=\mathrm{N}) ; \delta_{\mathrm{H}}\left(\mathrm{CDCl}_{3} ; 300 \mathrm{MHz}\right) 2.07$ and $2.28(2 \mathrm{~s}, 6 \mathrm{H}, \mathrm{Me}), 3.14$ $(\mathrm{d}, 3 \mathrm{H}, \mathrm{Me}, J=1.5 \mathrm{~Hz}), 3.75(\mathrm{~d}, 1 \mathrm{H}, \mathrm{CH}, J=8.3 \mathrm{~Hz}), 6.67$ (dq, $1 \mathrm{H}, \mathrm{CH}, J=1.5$ and 8.3 Hz ) and $7.1-7.5(\mathrm{~m}, 8 \mathrm{H}, \mathrm{Ar})$; $\delta_{\mathrm{C}}\left(\mathrm{CDCl}_{3} ; 75 \mathrm{MHz}\right) 17.54$ and $18.25(2 \mathrm{Me}), 45.37$ and 44.61 ( 2 C), $47.27(\mathrm{MeN}), 59.78(\mathrm{CH}), 118.86,119.06,122.13,123.24$, $126.04,126.18,126.64,126.85,141.12,143.83,145.42$ and 145.70 (Ar) and $163.69(\mathrm{C}=\mathrm{N})$ (HRMS Calcd. $\mathrm{C}_{19} \mathrm{H}_{19} \mathrm{NS}$; 293.1238. Found: 293.1230).

## 12-( $N$-Isopropyliminomethyl)-9,10,dihydro-11-thia-9,10-ethanoanthracene (4b)

A solution of isopropylamine ( 1 mmol ) in dichloromethane $\left(5 \mathrm{~cm}^{3}\right.$ ) containing molecular sieves was cooled to $0^{\circ} \mathrm{C}$. 12-Formyl-9,10-dimethyl-9,10-dihydro-11-thia-9,10-ethanoanthracene ( 1 mmol ) dissolved in dichloromethane $\left(5 \mathrm{~cm}^{3}\right)$ was added dropwise and the reaction was carried out for 8 h . After filtration and evaporation of the solvent, the desired imine was obtained in quantitative yield. No aldehyde was detected in its NMR and IR spectra. Mp $116-121^{\circ} \mathrm{C}$ (decomp.); $v\left(\mathrm{CDCl}_{3}\right) /$ $\mathrm{cm}^{-1} 1650(\mathrm{C}=\mathrm{N}) ; \delta_{\mathrm{H}}\left(\mathrm{CDCl}_{3} ; 300 \mathrm{MHz}\right) 1.02$ and $1.16(2 \mathrm{~d}, 6 \mathrm{H}$, 2 Me -isopropyl), 2.08 and $2.28(2 \mathrm{~s}, 6 \mathrm{H}, \mathrm{Me}), 3.15$ (m, $1 \mathrm{H}, \mathrm{CH}$ isopropyl), $3.73(\mathrm{~d}, 1 \mathrm{H}, \mathrm{CH}), 6.64(\mathrm{~d}, 1 \mathrm{H}, \mathrm{CH})$ and $7.1-7.4(\mathrm{~m}, 8$ $\mathrm{H}, \mathrm{Ar}) ; \delta_{\mathrm{C}}\left(\mathrm{CDCl}_{3}, 75 \mathrm{MHz}\right) 17.34$ and $17.86(2 \mathrm{Me}), 23.61$ and 24.33 ( 2 Me -isopropyl), 45.37 and 44.61 (2 C), 59.33 and 60.63 ( 2 CH ), $118.68,118.89,122.97,125.82,126.00,126.42,126.65$, $141.08,143.73,145.25$ and 145.64 (Ar) and $159.76(\mathrm{C}=\mathrm{N})$ (HRMS Calcd. $\mathrm{C}_{21} \mathrm{H}_{23} \mathrm{NS}, 321.1551$ Found: 321.1550 ).

## FVT of compounds $4 a$ and $b$

50 mg samples of imine $4 \mathbf{a}$ or $\mathbf{b}$ were sublimed under $10^{-5} \mathrm{hPa}$ through an empty quartz tube $(l=20 \mathrm{~cm}$, i.d. $=1.6 \mathrm{~cm})$ heated at $500^{\circ} \mathrm{C}$ (the FVT apparatus has been described elsewhere ${ }^{4.6 .22}$ ). DMA was separated by crystallization at the oven exit. The other products were condensed at $-196^{\circ} \mathrm{C}$ on a frozen mixture of $\mathrm{CDCl}_{3}$ and $\mathrm{CFCl}_{3}(3: 1)$, transferred upon melting in a NMR tube fitted below the cold trap, and immediately analysed by NMR spectroscopy at $-60^{\circ} \mathrm{C}$
(Bruker 250 apparatus), then at room temperature. The NMR spectra of the obtained dihydrothiazoles are reported in Scheme 3. IR data were reported in the preliminary communication. ${ }^{4}$

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[^0]:    $\dagger \mathrm{I} \mathrm{kcal} \mathrm{mol}^{-1}=4.184 \mathrm{~kJ} \mathrm{~mol}^{-1}$

